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# Merriam- Webster's Collegiate® Dictionary

TENTH EDITION

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worth two shillings b: any of several similar coins issued in Commonwealth countries 3: GULDEN 4: FORINT  
**florist** \flor-ist, -'flor-, 'flär- n (1623): one who sells or grows for sale flowers and ornamental plants — **florist-ry** \-s-trē n  
**floristic** \flō-ris-tik/ adj (1898): of or relating to flowers, a flora, or the phytogeographical study of plants and plant groups — **floristic-ally** \-ti-k(ə)-lē adv

**flourish** \flōr-(y)-wət, 'flär- n [L. be flourished, fr. *florere* to flourish] (1843): a period of flourishing (as of a person or movement)  
**floss** \flās, 'flōs/ n [prob. modif. of *Flos* soft, weak (of silk fiber), fr. Gascon, fr. L. *fluxus*, lit., loose, flowing, pp. of *fluere* to flow — more at **FLUID**] (1759) 1 a: soft thread of silk or mercerized cotton for embroidery b: DENTAL FLOSS 2: fluffy fibrous material  
**floss vt** (1974): to use dental floss on ~ vi: to use dental floss

**flossy** \flā-sē, 'flō- adj floss-i-er-, -est (1839) 1: of, relating to, or having the characteristics of floss 2: stylish or glamorous esp. at first impression (~ new hotels) — **floss-ily** \flā-sē-lē adv  
**float** \flō-tā/ n [Sp] (1527): a fleet of Spanish ships  
**floatation** \flō-tā-shən/ n [float] (1806) 1: the act, process, or state of floating 2: an act or instance of financing (as an issue of stock) 3: the separation of the particles of a mass of pulverized ore according to their relative capacity for floating on a given liquid; also: any of various similar processes involving the relative capacity of materials for floating 4: the ability (as of a tire or snowshoes) to stay on the surface of soft ground or snow

**float-la** \flō-ti-lā/ n [Sp, dim. of *flota* fleet, fr. OF *flote*, fr. ON *floti*; akin to OE *flota* ship, fleet — more at **FLAOT**] (1711) 1: a fleet of ships or boats; esp: a navy organizational unit consisting of two or more squadrons of small warships 2: an indefinite large number (a ~ of changes)  
**float-sam** \flāt-səm/ n [AF *floteson*, fr. OF *floter* to float, of Gmc origin; akin to OE *flotan* to float, *flota* ship] (ca. 1607) 1: floating wreckage of a ship or its cargo; broadly: floating debris 2 a: a floating population (as of emigrants or castaways) b: an accumulation of miscellaneous or unimportant stuff

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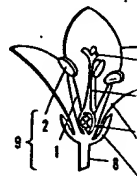
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ter esp. of a stream 2: gradual deformation of a body of (as rock) by intermolecular shear  
**flow-chart** \flō-čārt/ n (1920): a diagram that shows step progression through a procedure or system esp. using connect a set of conventional symbols — **flow-chart-ing** \-čārt-iŋ/ n  
**flow cy-to-m-e-try** \flō-si-tā-mē-trē/ n (1978): a technique ing and sorting cells and their components (as DNA) by a fluorescent dye and detecting the fluorescence usu. by illumination  
**flow diagram** n (1943): FLOWCHART  
**flower** \flāw-(ə)-r/ n [ME *flour* flower, best of anything, flour, fr. OF *flor*, flour, fr. L. *flor*, *flor* — more at **BLOW**] (13c) 1 a: BLOSSOM, INFLORESCENCE b: a shoot of the sporophyte of a higher plant that is modified for reproduction and consists of a shortened axis bearing modified leaves; esp: one of a seed plant differentiated into a calyx, corolla, stamens, and carpels c: a plant cultivated for its blossoms 2 a: the best part or example (the ~ of our youth) b: the finest most vigorous period c: a state of blooming or flourishing (in full ~) 3 pl: a finely divided powder produced esp. by condensation or sublimation (~s of sulfur) — **flowered** \flāw-(ə)-r-əd/ adj — **flower-ful** \flāw-(ə)-fəl/ adj — **flower-er-less** \flāw-(ə)-r-ləs/ adj — **flower-like** \-līk/ adj

**flower vi** (13c) 1 a: DEVELOP (~ed into young woman) b: FLOURISH 2: to produce flowers: BLOSSOM ~ vt 1: to bear flowers 2: to decorate with flowers or floral designs — **flower** \flāw-(ə)-r-ən/ n  
**flower-age** \flāw-(ə)-r-ij/ n (1840): a flowering process, station  
**flower bud** n (1828): a plant bud that produces only a flower  
**flower bug** n (ca. 1889): any of various small mostly black predaceous bugs (family Anthicidae) that frequent flowers on pest insects (as aphids and thrips)  
**flower child** n (1967): a hippie who advocates love, beauty, flower-er also flower-ette \flāw-(ə)-r-ət/ n (15c): FLORET 1: flower girl n (1925): a little girl who carries flowers at a wedding  
**flower head** n (1845): a capitulum (as of a composite) having flowers so arranged that the whole inflorescence looks like a flower  
**flowering dogwood** n (1843): a common spring-flowering bracted dogwood (*Cornus florida*)  
**flowering plant** n (1745): ANGIOSPERM  
**flower people** n pl (1967): FLOWER CHILDREN  
**flower-pot** \flāw-(ə)-pāt/ n (1598): a pot in which to grow  
**flower-ery** \flāw-(ə)-r-ē/ adj (14c) 1: of, relating to, or resembling 2: marked by or given to rhetorical elegance — **flower-er-ly** \flāw-(ə)-r-lē/ adv — **flower-i-ness** n  
**flower-meter** \flō-mē-tər/ n (1915): an instrument for measuring more properties (as velocity or pressure) of a flow (as of a pipe)  
**flown** \flōn/ past part of FLY  
**flown adj** [archaic pp. of *flow*] (1626): filled to excess  
**flow sheet** n (1912): FLOWCHART  
**flow-stone** \flō-s-tōn/ n (1925): calcite deposited by a thin flowing water usu. along the walls or floor of a cave  
**flu** \flū/ n [by shortening] (1839) 1: INFLUENZA 2: any virus diseases marked esp. by respiratory symptoms  
**flub** \flab/ vb flubbed; flubbing [origin unknown] vt (18) make a mess of: BOTCH (flubbed my lines) ~ vi: BLUNDER  
**flub** n (1948): an act or instance of flubbing  
**flub-dub** \flab-dəb/ n [origin unknown] (1888): BUNKUM, DASH

**fluc-tu-ant** \flak-čə-wənt/ adj (1560) 1: moving in waves: ABLE, UNSTABLE 3: being movable and compressible (a ~able fluc-tu-ate \flak-čə-wāt/ vb -at-ed, -at-ing [L. *fluctuare*, fr. *fluctus* flow, wave, fr. *fluere* — more at **FLUID**)] ~ vt: to shift back and forth uncertainly 2: to ebb and flow in ~ vi: to cause to fluctuate ~ syn see SWING — **fluc-tu-a-tion** \flū-šən/ n — **fluc-tu-a-tion-al** \flū-šən-əl, -shən-əl/ adj  
**flue** \flū/ n [origin unknown] (1582): an enclosed passage directing a current: as a: a channel in a chimney for conveying smoke to the outer air b: a pipe for conveying flame gases around or through water in a steam boiler c: an air leading to the lip of a wind instrument d: FLUE PIPE  
**flue-cured** \flū-kyurd/ adj (1905): cured with heat transmitted a flue without exposure to smoke or fumes (~ tobacco)  
**flu-en-cy** \flū-ən-sē/ n (1636): the quality or state of being **flu-ent** \flū-ənt/ adj [L. *fluens*, pp. of *fluere*] (1599) 1: capable of flowing: FLUID b: capable of moving with ease (~ the ~ body of a dancer) 2 a: ready or facile in speech (~ ish) b: effortlessly smooth and rapid: POLISHED (a ~ performer) — **flu-ent-ly** adv

**flue pipe** n (1852): an organ pipe whose tone is produced by current striking the lip and causing the air within to vibrate — **flue stop** n (1855): an organ stop made up of flue pipes  
**fluff** \flaf/ n [perh. blend of *fluff* (fluff) and *pufl*] (1790) 1: DOWN 2: something fluffy 3: something inconsequential 4: DER; esp: an actor's lapse of memory  
**fluff vi** (1875) 1: to become fluffy 2: to make a mistake, forget or bungle one's lines in a play ~ vt 1: to make fluffy: to spoil by a mistake: BOTCH b: to deliver badly or (in lines) in a play  
**fluffy** \flā-fē/ adj fluff-i-er-, -est (ca. 1825) 1 a: covered with fluffy b: being light and soft or airy (a ~ coat)



cross section of fl  
 filament, 2 anther,  
 style, 5 petal, 6 ovary,  
 8 pedicel, 7 stamen,  
 11 perianth.

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lacking in meaning or substance : SUPERFICIAL 2c ~ fluff-ily \ˈflʌ-  
 -li-ə-adv ~ fluff-i-ness \ˈflʌ-ɪ-ˈnəs-ŋ n  
 flū-gel-horn or flue-gel-horn \ˈflū-ɡəl-ˌhɔrn. ˈflū-ə- n [G, fr. *Flügel*  
 wing, flank + *Horn* horn, fr. its use to signal the flanking drivers in a  
 battle] (1854) : a valved brass instrument resembling a cornet but  
 having a larger bore ~ flū-gel-horn-ist \-hɔrn-ˈnist- n  
 fluid \ˈflū-ɪ-d-ə- adj [F or L; F *fluide*, fr. L *fluidus*, fr. *fluere* to flow;  
 akin to Gk *phlyzein* to boil over] (1603) 1 a : having particles that  
 easily move and change their relative position without a separation of  
 the mass and that easily yield to pressure : capable of flowing b  
 : subject to change or movement (boundaries became ~) 2 : charac-  
 terized by or employing a smooth easy style (the ballerina's ~ move-  
 ments) 3 a : available for a different use b : LIQUID 4 (~ assets) ~  
 -fluid-ly adv ~ fluid-ness n  
 fluid n (1661) : a substance (as a liquid or gas) tending to flow or con-  
 form to the outline of its container ~ fluid-al \ˈflū-ɪ-d-əl- adj ~ fluid-  
 -ly \-d-əl- adv  
 fluid dram or flu-id-dram \ˈflū-ɪ-d-ə-ˈdram- n (ca. 1860) : a unit of  
 liquid capacity equal to 1/8 fluid ounce — see WEIGHT table  
 fluid-ex-tract \ˈflū-ɪ-d-ə-ˈdek-ˌstrakt- n (1851) : an alcohol preparation of  
 a vegetable drug containing the active constituents of one gram of the  
 dry drug in each milliliter  
 fluid-ice \ˈflū-ɪ-d-ɪk- adj (1960) : of, relating to, or being a device (as an  
 amplifier or control) that depends for operation on the pressures and  
 flows of a fluid in precisely shaped channels ~ fluidic n ~ fluid-ics  
 \-d-ɪks- n pl but sing in constr  
 fluid-ity \ˈflū-ɪ-d-ə-ˈtɪ- n (1603) 1 : the quality or state of being fluid 2  
 : the physical property of a substance that enables it to flow  
 fluid-ize \ˈflū-ɪ-d-ə-ˈdɪz- v -ized; -izing- (ca. 1855) 1 : to cause to flow  
 like a fluid 2 : to suspend (as solid particles) in a rapidly moving  
 stream of gas or vapor to induce flowing motion of the whole ~ fluid-  
 -ization \ˈflū-ɪ-d-ə-ˈzə-ˌʃən- n ~ fluid-izer \ˈflū-ɪ-d-ə-ˈdɪ-zər- n  
 fluidized bed n (1949) : a bed of small solid particles (as in a coal  
 burning furnace) suspended and kept in motion by an upward flow of a  
 fluid (as a gas) — called also *fluid bed*  
 fluid mechanics n pl but sing or pl in constr (1937) : a branch of me-  
 chanics dealing with the properties of liquids and gases  
 fluid ounce n (ca. 1860) 1 : a U.S. unit of liquid capacity equal to 1/16  
 pint — see WEIGHT table 2 : a British unit of liquid capacity equal to  
 1/8 pint — see WEIGHT table  
 fluke \ˈflʌk- n [ME, fr. OE *flōc*; akin to OE *flōh* chip, OHG *flah*  
 smooth, Gk *plax* flat surface, and prob. to OE *flōr* floor — more at  
 FLOOR] (bef. 12c) 1 : FLATFISH 2 : a flattened digenetic trematode  
 worm; broadly : TREMATODE ~ compare LIVER FLUKE  
 fluke n [perh. fr. */fluke/*] (1561) 1 : the part of an anchor that fastens  
 in the ground — see ANCHOR illustration 2 : one of the lobes of a  
 whale's tail  
 fluke n [origin unknown] (1857) 1 : an accidentally successful stroke  
 at billiards or pool 2 : a stroke of luck (the discovery was a ~)  
 fluky also fluk-ey \ˈflū-ke-ə- adj fluk-i-er; -est (1867) 1 : happening  
 by or depending on chance 2 : being unsteady or uncertain — used  
 esp. of wind  
 flume \ˈflʌm- n [prob. fr. ME *flum* river, fr. OF, fr. L *flumen*, fr. *fluere*  
 — more at FLUID] (1748) 1 : an inclined channel for conveying water  
 (as for power) 2 : a ravine or gorge with a stream running through it  
 flum-mery \ˈflʌm-ri-, ˈflʌ-mə- n, pl -mer-ies [W *lylrum*] (1623) 1 a  
 : a soft jelly or porridge made with flour or meal b : any of several  
 sweet desserts 2 : MUMMERY, MUMBO JUMBO  
 flum-mox \ˈflʌ-mɔks-, -mɪks- v [origin unknown] (1837) : CONFUSE  
 flump \ˈflʌmp- [imit.] v (1816) : to move or fall suddenly and heavily  
 (~ed down into the chair) ~ vt : to place or drop with a flump  
 flump n (1832) : a dull heavy sound (as of a fall)  
 flung past and past part of FLING  
 flunk \ˈflʌŋk- vb [perh. blend of *flinch* and *funk*] vi (1823) : to fail esp.  
 in an examination or course ~ vt 1 : to give a failing grade to 2 : to  
 get a failing grade in ~ flunk-er n  
 flunk n (1846) : an act or instance of flunking  
 flunk out vi (1920) : to be dismissed from a school or college for failure  
 ~ vt : to dismiss from a school or college for failure  
 flunk-y or flun-key \ˈflʌŋ-ke-ə- n, pl flunkies or flunkies [Sc, of un-  
 known origin] (ca. 1782) 1 a : a livered servant b : one performing  
 menial or miscellaneous duties 2 : YES-MAN  
 fluo-cin-o-lone ac-e-to-nide \ˈflū-ɔ-ˌsi-ˌn-ˌɪ-, ˌɒn-ə-ˌsə-ˈ(ɪ)-ˌnɪd- n [fluor-  
 + *cinolone*, prob. alter. of *nisolone* (as in *prednisolone*)] (1963) : a  
 glucocorticoid steroid C<sub>21</sub>H<sub>27</sub>F<sub>5</sub>O<sub>6</sub> used esp. as an anti-inflammatory  
 agent in the treatment of skin diseases  
 fluor or flū-ōr, ˈflū-ōr- n [NL, mineral belonging to a group used as  
 fluxes and including fluoric, fr. L, flow, fr. *fluere* — more at FLUID]  
 (1661) : FLUORITE  
 fluor- or fluoro- comb form [F] 1 : fluorine (<fluoride). 2 also fluori-  
 : fluorescence (<fluoresce) (<fluorimeter)  
 fluo-resce \ˈflū-ˌres-, ˌflō- v -resced; -resc-ing [back-formation fr.  
 fluorescence] (1874) : to produce, undergo, or exhibit fluorescence ~  
 fluo-res-cer n  
 fluo-res-cen-tial \ˈre-s-ə-ˌn- n (1876) : a yellow or red crystalline dye  
 CaF<sub>2</sub>O<sub>3</sub> with a bright yellow-green fluorescence in alkaline solution  
 fluo-res-cen-tial \ˈre-s-ə-ˌn- n (<fluorspar + opalescence) (1852) : lumi-  
 nescence that is caused by the absorption of radiation at one wave-  
 length followed by nearly immediate reradiation usu. at a different  
 wavelength and that ceases almost immediately when the incident  
 radiation stops; also : the radiation emitted — compare PHOSPHORESC-  
 -ence  
 fluo-res-cent \-s-nt- adj (1853) 1 : having or relating to fluorescence  
 2 : bright and glowing as a result of fluorescence (~ inks); broadly  
 : very bright in color ~ fluorescent n  
 fluorescent lamp n (1896) : a tubular electric lamp having a coating  
 of fluorescent material on its inner surface and containing mercury  
 vapor, whose bombardment by electrons from the cathode provides  
 ultraviolet light which causes the material to emit visible light  
 fluo-r-date \ˈflū-ɪ-d-ət, ˈflōr-, ˈflōr- v -dat-ed; -dat-ing (1949) : to add  
 a fluoride to (as drinking water) to reduce tooth decay ~ fluo-ri-  
 -dation \ˈflū-ɪ-d-ə-ˌʃən-, ˈflōr-, ˈflōr- n  
 fluoride \ˈflōr-ɪd-, ˈflū-ɪ- n, often attrib (1826) 1 : a compound of  
 fluorine 2 : the monovalent anion of fluorine

fluor-*n*ate \flōr-ə-nāt, flōr-, flūr-, *n* -nat-ed; -nat-ing (ca. 1929) : to treat or cause to combine with fluorine or a compound of fluorine / — fluo-*r*-i-na-tion \flōr-ə-nā-shən, flōr-, flūr-, *n* fluo-*r*-ine \flūr-ēn, flōr-, flōr-*n* [F, fr. NL *fluor*] (1813) : a nonmetallic halogen element that is isolated as a pale yellowish flammable irritating toxic diatomic gas — see ELEMENT table  
fluor-ite \flūr-īt, flōr-, flōr-*n* [It, fr. NL *fluor*] (1868) : a transparent or translucent mineral of different colors that consists of the fluoride of calcium and is used as a flux and in the making of opalescent and opaque glasses  
fluor-o-car-bon \flūr-ə-kär-bən, flōr-, flōr-*n* (1937) : any of various chemically inert compounds containing carbon and fluorine used chiefly as lubricants, refrigerants, nonstick coatings, and formerly aerosol propellants and in n.aking resins and plastics; also : CHLOROFLUOROCARBON  
fluor-o-chrome \flūr-ə-krōm, flōr-, flōr-*n* (1943) : any of various fluorescent substances used in biological staining to produce fluorescence in a specimen  
fluor-ro-graphy \flūr-rä-grä-fē, flō-, flōr-, *n* (1941) : PHOTOFUOROGRAPHY — fluo-*r*-o-graph-ic \flūr-ə-grä-fik, flōr-, flōr-*adj*  
fluor-rom-e-ter \flūr-rä-mä-tēr, flō-, flō- or fluo-rim-e-ter \-ri-*n* (1897) : an instrument for measuring fluorescence and related phenomena (as intensity of radiation) — fluo-*r*-met-ric \flūr-rä-mē-trik, flō-, flō- or fluo-rim-e-try \flūr-rä-mē-trē, flō-, flō- or fluo-rim-e-try \-ri-mē-trē *n*  
fluor-ro-scope \flūr-rä-skōp, flōr-, flōr-*n* [ISV] (1896) : an instrument used for observing the internal structure of an opaque object (as the living body) by means of X rays — fluo-*r*-ro-scop-ic \flūr-ə-skä-pik, flōr-, flōr-*adj* — fluo-*r*-ro-scop-i-cal-ly \-pi-k(ə)-lēv *adv* — fluo-ros-cop-ist \flūr-räs-kä-pist, flō-, flō-*n* — fluo-ros-cop-y \-pē *n*  
fluoroscope *vi* -scop-ed; -scop-ing (1898) : to examine by fluoroscopy  
fluor-ro-sis \flūr-rä-sas, flō-*n* [NL] (1927) : an abnormal condition (as mottling of the teeth) caused by fluorine or its compounds — fluo-rot-ic \-rä-tik *adj*  
fluor-u-ura-cil \flūr-ə-yūr-ə-sil, -səl, flōr-, flōr-*n* [fluor- + uracil] (ca. 1958) : a fluorine-containing pyrimidine base C<sub>4</sub>H<sub>3</sub>FN<sub>2</sub>O<sub>2</sub> used to treat some kinds of cancer  
fluor-spar \flūr-spär, flōr-, flōr-*n* (1794) : FLUORITE  
flu-phen-azine \flūr-fē-nä-zēn \ *n* [fluor- + phenazine] (ca. 1960) : a tranquilizing compound C<sub>12</sub>H<sub>8</sub>F<sub>3</sub>N<sub>3</sub>O<sub>2</sub> used esp. combined as a salt  
flur-ry \flūr-ē, flā-rē *n*, pl flurr-ies [prob. fr. *flurr* (to throw scatteringly)] (1686) 1 *a* : a gust of wind 2 *a* : a brief light snowfall 2 *a* : a brief period of commotion or excitement 3 *a* : sudden occurrence of many things at once : BARRAGE 2 (as of insults) 3 *a* : brief advance or decline in prices : a short-lived outburst of trading activity  
flurry *vb* flur-ried; flur-ry-ing *vi* (1757) : to cause to become agitated and confused ~ *vi* : to move in an agitated or confused manner  
flush \fləsh \ *vb* [ME *flusshen*] *vi* (13c) : to take wing suddenly ~ *vi* 1 : to cause (a bird) to flush 2 : to expose or chase from a place of concealment (~ed the boys from their hiding place)  
flush *n* [MF *flus*, *fluz*, fr. L *fluxus* flow, flux] (ca. 1529) 1 : a hand of playing cards all of the same suit; *specif* : a poker hand containing five cards of the same suit but not in sequence — see POKER illustration 2 : a series of three or more slalom gates set vertically on a slope  
flush *n* [perh. modif. of L *fluxus*] (1529) 1 : a sudden flow (as of water); also : a rinsing or cleansing with or as if with a flush of water 2 *a* : a sudden increase or expansion; esp : sudden and usu. abundant new plant growth 3 *a* : surge of emotion (felt a ~ of anger at the insult) 3 *a* : a tinge of red : BLUSH 3 *a* : fresh and vigorous state (in the first ~ of womanhood) 4 : a transitory sensation of extreme heat — compare HOT FLASH  
flush *vi* (1548) 1 : to flow and spread suddenly and freely 2 *a* : to glow brightly 3 *a* : BLUSH 3 : to produce new growth (the plants ~ twice during the year) *vi* 1 *a* : to cause to flow 2 *a* : to pour liquid over or through; esp : to cleanse or wash out with or as if with a rush of liquid (~ the toilet) (~ the lungs with air) 2 : INFLAME, EXCITE — usu. used passively (~ed with pride) 3 : to cause to blush  
flush *acc*, (1594) 1 *a* : of a ruddy healthy color 3 *a* : full of life and vigor : LUSTY 2 *a* : filled to overflowing 3 *a* : AFFLUENT 3 : readily available : ABUNDANT 4 *a* : having or forming a continuous plane or unbroken surface (~ panning) 4 *a* : directly abutting or immediately adjacent : as (1) : set even with an edge of a type page or column : having no indentation (2) : arranged edge to edge so as to fit snugly — flush-ness *n*  
flush *adv* (1700) 1 : in a flush manner 2 : SQUARELY (hit him ~ on the chin)  
flush *vi* (ca. 1842) : to make flush (~ the headings on a page)  
flush-able \flā-shə-bəl *adj* (1973) : suitable for disposal by flushing down a toilet  
flush-ter \flāst-ər *vb* flush-tered; flush-ter-ing \-(tə-)rīg [prob. of Scand origin; akin to Icel *flástur* hurry] *vi* (1604) 1 : to make tipsy 2 : to put into a state of agitated confusion : UPSET ~ *vi* : to move or behave in an agitated or confused manner *syn* see DISCOMPOSE — flush-tered-ly *adv*  
fluster *n* (1728) : a state of agitated confusion  
flute \flüt-*n* [ME *floute*; fr. MF *flaute*, fr. OF *flaüte*, prob. fr. OProv *flaut*] (14c) 1 *a* : RECORDER 3 *a* : keyed woodwind instrument consisting of a cylindrical tube which is stopped at one end and which has a side hole over which air is blown to produce the tone and having a range from middle C upward for three octaves 2 : something



ʌ\ about ʌ\ kitten, F table ʌr\ further ʌ\ ash ʌ\ ace ʌ\ mop, mar  
 ʌu\ out ʌh\ chin ʌt\ bet ʌl\ easy ʌg\ go ʌi\ hit ʌ\ ice ʌj\ job  
 ʌj\ sing ʌl\ go ʌw\ law ʌoi\ boy ʌth\ thing ʌh\ the ʌü\ lagoon ʌü\ foot  
 ʌy\ yet ʌzh\ vision ʌ, ɛ, ʰ, æ, œ, u, ʉ, ʰ See Guide to Pronunciation

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# KINETIC THEORY OF GASES AND THE BOLTZMANN EQUATION

In Section 7-2 we introduced the concept of phase space and distribution functions in phase space. We also derived the Liouville equation, which is the equation of motion that the phase space distribution function must satisfy. Since we were interested only in equilibrium statistical mechanics at that time, we did not consider the Liouville equation in any detail. In this chapter we shall review the concept of phase space and derive the Liouville equation again. We shall then introduce reduced distribution functions and derive the Bogoliubov, Born, Green, Kirkwood, Yvon (BBGKY) hierarchy. This hierarchy is the nonequilibrium generalization of the Kirkwood integral equation hierarchy for the fluid distribution functions,  $g^{(n)}(\mathbf{r}_1, \dots, \mathbf{r}_n)$ , of Chapter 13. Nobody has yet devised a successful way to uncouple the BBGKY hierarchy, and so in Section 18-4 we shall derive a physical, yet approximate, equation for the distribution function for gases. This equation, called the Boltzmann equation, is the central equation of the rigorous kinetic theory of gases. In Section 18-5, we shall derive some of the general consequences of the Boltzmann equation that can be determined without actually solving it completely. We shall discuss its solution in Chapter 19. The standard reference for most of this chapter is Hirschfelder, Curtiss, and Bird. Mazo (in "Additional Reading") also discusses these topics well.

## 18-1 PHASE SPACE AND THE LIOUVILLE EQUATION

Consider a system of  $N$  point particles. The classical dynamical state of this system is specified by the  $3N$  momentum components  $p_1, p_2, \dots, p_{3N}$  and the  $3N$  spatial coordinates  $q_1, \dots, q_{3N}$ . We can construct a  $6N$ -dimensional space whose coordinates are  $q_1, q_2, \dots, p_1, \dots, p_{3N}$ . One point in this *phase space* completely specifies the microscopic dynamical state of our  $N$ -particle system. As the system evolves in time, this *phase point* moves through phase space in a manner completely dictated by the equations of motion of the system. Actually, one never knows (nor really cares to know) the  $6N$  coordinates of a macroscopic system. Rather, one knows just a few

macroscopic mechanical properties of the system, such as the energy, volume, velocity, etc. Clearly there are a great number of points in phase space that are compatible with the few variables that we know about the system. The set of all such phase points constitutes an *ensemble* of systems. The number of systems in an ensemble approaches infinity, and so the set of phase points that could possibly represent our system becomes quite dense. This allows us to define a *density of phase points* or *distribution function* as the fraction of phase points contained in the volume  $dq_1 dq_2 \cdots dp_{3N}$ . We shall denote the phase space distribution function by  $f_N(q_1, q_2, \dots, p_{3N}, t)$ , or more conveniently by  $f_N(p, q, t)$ . We shall often use this abbreviated notation. Similarly, we shall often denote  $dq_1, dq_2 \cdots dp_{3N}$  by  $dp dq$ . The density  $f_N(p, q, t)$  is normalized such that

$$\int f_N(p, q, t) dp dq = 1$$

Since each phase point moves in time according to the equations of motion of the system it describes,  $f_N$  itself must obey some sort of equation of motion. The equation that  $f_N(p, q, t)$  satisfies can be readily determined by using the methods of the previous chapter, particularly, the argument associated with Eqs. (17-1) to (17-5). The number of phase points within some arbitrary volume  $v$  is

$$n = \mathcal{N} \int_v f_N(p, q, t) dp dq$$

where we are using the condensed notation of letting  $p$  and  $q$  denote all the spatial coordinates and momenta necessary to specify a system in the ensemble. The rate of change of the number of phase points within  $v$  is

$$\frac{dn}{dt} = \mathcal{N} \int_v \frac{\partial f_N}{\partial t} dp dq \quad (18-1)$$

Since phase points are neither created nor destroyed, the rate of change of  $n$  must be given by the rate at which phase points flow through the surface enclosing  $v$ . The rate of flow of phase points is  $\mathcal{N} f_N \mathbf{u}$ , where  $\mathbf{u}$  is not just the  $3N$ -dimensional vector  $(\dot{q}_1, \dot{q}_2, \dots, \dot{q}_{3N})$ , but the  $6N$ -dimensional vector  $(\dot{q}_1, \dots, \dot{p}_1, \dots, \dot{p}_{3N})$  since the spatial coordinates and momenta play an equivalent role in phase space. We integrate this flow over the surface to get

$$\frac{dn}{dt} = -\mathcal{N} \int_S f_N \mathbf{u} \cdot d\mathbf{S}$$

The negative sign here indicates that an outflow of phase points yields a negative value for  $dn/dt$  since  $\mathbf{u} \cdot d\mathbf{S}$  is positive if  $\mathbf{u}$  is directed outward from  $v$  and negative if  $\mathbf{u}$  is directed inward.

The surface integral can be transformed to a volume integral by using Gauss' theorem to get

$$\frac{dn}{dt} = -\mathcal{N} \int_v \nabla \cdot (f_N \mathbf{u}) dp dq \quad (18-2)$$

If we subtract Eq. (18-1) from Eq. (18-2) and realize that this equation is valid for any choice of  $v$ , we have the equation for the conservation of phase points

$$\frac{\partial f_N}{\partial t} + \nabla \cdot (f_N \mathbf{u}) = 0 \quad (18-3)$$



in which it should be clear that since we are dealing with phase space

$$\mathbf{u} = (\dot{q}_1, \dots, \dot{q}_{3N}, \dot{p}_1, \dots, \dot{p}_{3N})$$

and

$$\begin{aligned} \nabla \cdot f_N \mathbf{u} &= \sum_{j=1}^{3N} \frac{\partial}{\partial q_j} (f_N \dot{q}_j) + \sum_{j=1}^{3N} \frac{\partial}{\partial p_j} (f_N \dot{p}_j) \\ &= \sum_{j=1}^{3N} \left\{ \frac{\partial f_N}{\partial q_j} \dot{q}_j + \frac{\partial f_N}{\partial p_j} \dot{p}_j \right\} + \sum_{j=1}^{3N} \left\{ \frac{\partial \dot{q}_j}{\partial q_j} + \frac{\partial \dot{p}_j}{\partial p_j} \right\} f_N \end{aligned}$$

But Eq. (7-27) shows that the summand of the second summation here is zero, and so Eq. (18-3) becomes

$$\frac{\partial f_N}{\partial t} + \sum_{j=1}^{3N} \frac{\partial f_N}{\partial q_j} \dot{q}_j + \sum_{j=1}^{3N} \frac{\partial f_N}{\partial p_j} \dot{p}_j = 0 \quad (18-4)$$

Using Hamilton's equations of motion,

$$\dot{p}_i = -\frac{\partial H}{\partial q_i} \quad \text{and} \quad \dot{q}_i = \frac{\partial H}{\partial p_i}$$

Eq. (18-4) can be written

$$\frac{\partial f_N}{\partial t} + \sum_{j=1}^{3N} \left( \frac{\partial H}{\partial p_j} \frac{\partial f_N}{\partial q_j} - \frac{\partial H}{\partial q_j} \frac{\partial f_N}{\partial p_j} \right) = 0 \quad (18-5)$$

The summation here is called a Poisson bracket and is commonly denoted by  $\{H, f_N\}$ ; so Eq. (18-5) is often written as

$$\frac{\partial f_N}{\partial t} + \{H, f_N\} = 0 \quad (18-6)$$

This is the Liouville equation, the most fundamental equation of statistical mechanics. In fact, it can be shown that the Liouville equation is equivalent to the  $6N$  Hamilton equations of motion of the  $N$ -body system.\*

In Cartesian coordinates, the Liouville equation reads

$$\frac{\partial f_N}{\partial t} + \sum_{j=1}^N \frac{\mathbf{p}_j}{m_j} \cdot \nabla_{\mathbf{r}_j} f_N + \sum_{j=1}^N \mathbf{F}_j \cdot \nabla_{\mathbf{p}_j} f_N = 0 \quad (18-6')$$

In this equation  $\nabla_{\mathbf{r}_j}$  denotes the gradient with respect to the spatial variables in  $f_N$ ;  $\nabla_{\mathbf{p}_j}$  denotes the gradient with respect to the momentum variables in  $f_N$ ; and  $\mathbf{F}_j$  is the total force on the  $j$ th particle.

One often sees the Liouville equation written as

$$i \frac{\partial f_N}{\partial t} = L f_N \quad (18-7)$$

where  $L$  is the Liouville operator,

$$L = -i \left( \sum_{j=1}^N \frac{\mathbf{p}_j}{m_j} \cdot \nabla_{\mathbf{r}_j} + \sum_{j=1}^N \mathbf{F}_j \cdot \nabla_{\mathbf{p}_j} \right) \quad (18-8)$$

\* Mazo, "Additional Reading," p. 23; M. Beran, *Amer. J. Phys.*, 35, p. 242, 1967.

The Liouville operator has been defined in such a way as to bring the Liouville equation into the form of the Schrödinger equation. A formal, and sometimes useful, solution to Eq. (18-7) is

$$f_N(\mathbf{p}, \mathbf{r}, t) = e^{-iLt} f_N(\mathbf{p}, \mathbf{r}, 0) \quad (18-9)$$

Note that the operator  $\exp(-iLt)$  displaces  $f_N$  ahead a distance  $t$  in time. This operator is called the time displacement operator of the system.

## 18-2 REDUCED DISTRIBUTION FUNCTIONS

Once we have the distribution function  $f_N(p, q, t)$ , we may compute the ensemble average of any dynamical variable,  $A(p, q, t)$ , from the equation

$$\langle A(t) \rangle = \int A(p, q, t) f_N(p, q, t) dp dq \quad (18-10)$$

It turns out that the dynamical variables of interest are functions of either the coordinates and momenta of just a few particles or can be written as a sum over such functions. A familiar example of this is the total intermolecular potential of the system. To a good approximation, this can be written as a sum over pair-wise potentials, and so

$$\langle U \rangle = \sum_{i,j} \int \cdots \int u(\mathbf{r}_i, \mathbf{r}_j) f_N(\mathbf{r}_1, \dots, \mathbf{p}_N, t) d\mathbf{r}_1 \cdots d\mathbf{p}_N \quad (18-11)$$

We encountered similar integrands when we studied the equilibrium theory of liquids. There we integrated over the coordinates of all the particles except  $i$  and  $j$  and called the resulting function of  $\mathbf{r}_i$  and  $\mathbf{r}_j$  a radial distribution function. We do the same thing here. We define reduced distribution functions  $f_N^{(n)}(\mathbf{r}_1, \dots, \mathbf{r}_n, \mathbf{p}_1, \dots, \mathbf{p}_n, t)$  by

$$f_N^{(n)}(\mathbf{r}_1, \dots, \mathbf{r}_n, \mathbf{p}_1, \dots, \mathbf{p}_n, t) = \frac{N!}{(N-n)!} \int \cdots \int f_N(\mathbf{r}_1, \dots, \mathbf{p}_N, t) d\mathbf{r}_{n+1} \cdots d\mathbf{r}_N d\mathbf{p}_{n+1} \cdots d\mathbf{p}_N \quad (18-12)$$

We shall usually drop the  $N$  subscript and furthermore write this simply as  $f^{(n)}(\mathbf{r}^n, \mathbf{p}^n, t)$ . Usually only  $f^{(1)}$  and  $f^{(2)}$  are necessary, and therefore we want to derive an equation for  $f^{(1)}$  and  $f^{(2)}$ . To do this, write the force  $\mathbf{F}_j$  appearing in the Liouville equation as the sum of the forces due to the other molecules in the system  $\sum_i \mathbf{F}_{ij}$  and an external force  $\mathbf{X}_j$ . Then multiply through by  $N!/(N-n)!$  and integrate over  $d\mathbf{r}_{n+1} \cdots d\mathbf{r}_N d\mathbf{p}_{n+1} \cdots d\mathbf{p}_N$  to get (Problem 18-2)

$$\begin{aligned} \frac{\partial f^{(n)}}{\partial t} + \sum_{j=1}^n \frac{\mathbf{p}_j}{m_j} \cdot \nabla_{\mathbf{r}_j} f^{(n)} + \sum_{j=1}^n \mathbf{X}_j \cdot \nabla_{\mathbf{p}_j} f^{(n)} \\ + \frac{N!}{(N-n)!} \sum_{i,j=1}^n \int \cdots \int \mathbf{F}_{ij} \cdot \nabla_{\mathbf{p}_j} f d\mathbf{r}_{n+1} \cdots d\mathbf{r}_N d\mathbf{p}_{n+1} \cdots d\mathbf{p}_N = 0 \end{aligned} \quad (18-13)$$

We have used the fact that  $f$  vanishes outside the walls of the container and when  $\mathbf{p}_i = \pm \infty$ . The last term in Eq. (18-13) can be broken up into two parts:

$$\begin{aligned} \sum_{i,j=1}^n \mathbf{F}_{ij} \cdot \nabla_{\mathbf{p}_j} f^{(n)} \\ + \frac{N!}{(N-n)!} \sum_{j=1}^n \sum_{i=n+1}^N \int \cdots \int \mathbf{F}_{ij} \cdot \nabla_{\mathbf{p}_j} f d\mathbf{r}_{n+1} \cdots d\mathbf{r}_N d\mathbf{p}_{n+1} \cdots d\mathbf{p}_N \end{aligned}$$

The second term here can be written as

$$\sum_{j=1}^n \iint \mathbf{F}_{j,n+1} \cdot \nabla_{\mathbf{p}_j} f^{(n+1)} d\mathbf{r}_{n+1} d\mathbf{p}_{n+1}$$

Putting all this together finally gives an exact equation for  $f^{(n)}$ , namely, (Problem 18-3),

$$\begin{aligned} \frac{\partial f^{(n)}}{\partial t} + \sum_{j=1}^n \frac{\mathbf{p}_j}{m_j} \cdot \nabla_{\mathbf{r}_j} f^{(n)} + \sum_{j=1}^n \mathbf{X}_j \cdot \nabla_{\mathbf{p}_j} f^{(n)} \\ + \sum_{i,j=1}^n \mathbf{F}_{ij} \cdot \nabla_{\mathbf{p}_j} f^{(n)} + \sum_{j=1}^n \iint \mathbf{F}_{j,n+1} \cdot \nabla_{\mathbf{p}_j} f^{(n+1)} d\mathbf{r}_{n+1} d\mathbf{p}_{n+1} = 0 \end{aligned} \quad (18-14)$$

This is the so-called Bogoliubov, Born, Green, Kirkwood, Yvon (BBGKY) hierarchy. This is the time-dependent generalization of the hierarchy that we derived earlier in the equilibrium theory of fluids. In fact, if one assumes that

$$f^{(n)} = g^{(n)}(\mathbf{r}_1, \dots, \mathbf{r}_n) \exp\left\{-\frac{1}{2mkT} \sum_{j=1}^n p_j^2\right\}$$

multiplies Eq. (18-14) through by  $p_i$ ,  $1 \leq i \leq n$ , and integrates over all momenta, one obtains the equilibrium hierarchy for  $g^{(n)}(\mathbf{r}_1, \dots, \mathbf{r}_n)$  (Problem 18-4). It would seem natural at this point to truncate this hierarchy by some sort of a superposition approximation, but so far this approach has not been successful.\* We shall end up deriving approximate equations for  $f^{(1)}$  and  $f^{(2)}$ .

Everything we have done up to now has been independent of density; i.e., it has been applicable to any density. Now we shall specialize to systems of dilute gases.

### 18-3 FLUXES IN DILUTE GASES

In a dilute gas, most of the molecules are not interacting with any other molecule and are just traveling along between collisions. Because of this, the macroscopic properties of a gas depend upon only the singlet distribution function  $f_j^{(1)}(\mathbf{r}, \mathbf{p}_j, t)$ . The subscript  $j$  here denotes the singlet distribution function of species  $j$ . This is the central distribution function of any theory of transport in dilute gases. In this section we shall define a number of averages over  $f_j^{(1)}$  and derive molecular expressions for the important flux quantities in terms of integrals over  $f_j^{(1)}$ . Since we shall be concerned only with gases in this and the following sections, we shall drop the superscript (1) from here on. We shall also write our equations in velocity space rather than momentum space, and so the distribution function of interest becomes  $f_j(\mathbf{r}, \mathbf{v}_j, t)$ . We shall renormalize  $f_j$  such that the integral of this distribution function over all velocities is the number density of  $j$  particles at the point  $\mathbf{r}$  at time  $t$ , i.e.,

$$\rho_j(\mathbf{r}, t) = \int f_j(\mathbf{r}, \mathbf{v}_j, t) d\mathbf{v}_j \quad (18-15)$$

Furthermore, if  $N_j$  is the total number of  $j$  molecules in our system, then

$$N_j = \iint f_j(\mathbf{r}, \mathbf{v}_j, t) d\mathbf{r} d\mathbf{v}_j \quad (18-16)$$

We shall now define a number of important average velocities.  $\mathbf{v}_j$  is the linear

\* See, for example, R. G. Mortimer, *J. Chem. Phys.*, 48, p. 1023, 1968.

velocity of a molecule of species  $j$ ; i.e., it is the velocity with respect to a coordinate system fixed in space. The average velocity is given by

$$\mathbf{v}_j(\mathbf{r}, t) = \frac{1}{\rho_j} \int \mathbf{v}_j f(\mathbf{r}, \mathbf{v}_j, t) d\mathbf{v}_j \quad (18-17)$$

and represents the macroscopic flow of species  $j$ . The *mass average velocity* is defined by

$$\mathbf{v}_0(\mathbf{r}, t) = \frac{\sum_j m_j \rho_j \mathbf{v}_j}{\sum_j m_j \rho_j} \quad (18-18)$$

Note that the denominator here is the mass density  $\rho_m(\mathbf{r}, t)$ . This velocity is often called the *flow velocity* or *stream velocity*. The momentum density of the gas is the same as if all the molecules were moving with velocity  $\mathbf{v}_0$ . The *peculiar velocity* is the velocity of a molecule relative to the flow velocity. The peculiar velocity  $\mathbf{V}_j$  is

$$\mathbf{V}_j = \mathbf{v}_j - \mathbf{v}_0 \quad (18-19)$$

The average of this peculiar velocity is the *diffusion velocity* (Problem 18-5). Clearly,

$$\bar{\mathbf{V}}_j = \frac{1}{\rho_j} \int (\mathbf{v}_j - \mathbf{v}_0) f_j(\mathbf{r}, \mathbf{v}_j, t) d\mathbf{v}_j \quad (18-20)$$

It is easy to show that (Problem 18-6)

$$\sum_j \rho_j m_j \bar{\mathbf{V}}_j = 0 \quad (18-21)$$

When we studied the elementary kinetic theory of gases, we saw that the various transport coefficients were related to molecular transport of mass, momentum, and kinetic energy. Let these molecular properties be designated collectively by  $\psi_j$ , where  $j$  refers to the particular species. We now derive expressions for the fluxes of these properties. Figure 18-1 shows a surface  $dS$  moving with velocity  $\mathbf{v}_0$ . The quantity  $\mathbf{n}$  is a unit vector normal to  $dS$ , and  $dS = \mathbf{n} dS$ . All the molecules that have velocity  $\mathbf{V}_j = \mathbf{v}_j - \mathbf{v}_0$  and that cross  $dS$  in the time interval  $(t, t + dt)$  must have been in a cylinder of length  $|\mathbf{V}_j| dt$  and base  $dS$ . This cylinder is shown in Fig. 18-1 and has a volume  $(\mathbf{n} \cdot \mathbf{V}_j) dS dt$ . Since there are  $f_j d\mathbf{v}_j$  molecules per unit volume with relative velocity  $\mathbf{V}_j$ , the number of  $j$  molecules that cross  $dS$  in  $dt$  is given by

$$(f_j d\mathbf{v}_j)(\mathbf{n} \cdot \mathbf{V}_j) dS dt$$

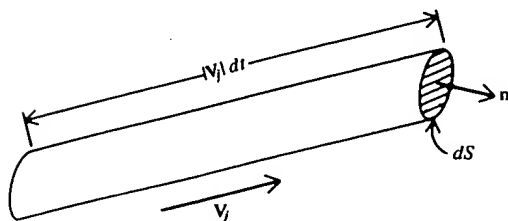


Figure 18-1. The cylinder containing all those molecules of species  $j$  with velocity  $\mathbf{V}_j$ , which cross the surface  $dS$  during the time interval  $dt$ . (From J. O. Hirschfelder, C. F. Curtiss, and R. B. Bird, *Molecular Theory of Gases and Liquids*. New York: Wiley, 1954.)

If each molecule carries with it a property  $\psi_j$ , then the flux of this property is

$$\psi_j f_j (\mathbf{n} \cdot \mathbf{V}_j) d\mathbf{v}_j$$

and the total flux across this surface is

$$\text{total flux} = \int \psi_j f_j (\mathbf{n} \cdot \mathbf{V}_j) d\mathbf{v}_j = \mathbf{n} \cdot \int \psi_j f_j \mathbf{V}_j d\mathbf{v}_j = \mathbf{n} \cdot \boldsymbol{\psi}_j \quad (18-22)$$

The vector  $\boldsymbol{\psi}_j$ ,

$$\boldsymbol{\psi}_j = \int \psi_j f_j \mathbf{V}_j d\mathbf{v}_j \quad (18-23)$$

is called the flux vector associated with the property  $\psi_j$ . The component of this vector in any direction is the transport of the property  $\psi_j$  in that direction. Let us now consider the various examples of  $\psi_j$ .

#### TRANSPORT OF MASS

In this case,  $\psi_j = m_j$ , and

$$\boldsymbol{\psi}_j = m_j \int f_j \mathbf{V}_j d\mathbf{v}_j = \rho_j m_j \overline{\mathbf{V}_j} \equiv \mathbf{j}_j \quad (18-24)$$

#### TRANSPORT OF MOMENTUM

Here  $\psi_j = m_j V_{jx}$ , and

$$\boldsymbol{\psi}_j = m_j \int V_{jx} f_j \mathbf{V}_j d\mathbf{v}_j = \rho_j m_j \overline{V_{jx} \mathbf{V}_j} \quad (18-25)$$

which is the flux of the  $x$ -component of momentum relative to  $\mathbf{v}_0$ . The flux of momentum is a pressure, which has components

$$(p_j)_{xx} = \rho_j m_j \overline{V_{jx} V_{jx}}$$

$$(p_j)_{xy} = \rho_j m_j \overline{V_{jx} V_{jy}}, \text{ etc}$$

or, in general,

$$p_j = \rho_j m_j \overline{\mathbf{V}_j \mathbf{V}_j} \quad (18-26)$$

which is the partial pressure tensor of the  $j$ th species.

#### TRANSPORT OF KINETIC ENERGY

$$\psi_j = \frac{1}{2} m_j V_j^2$$

and

$$\boldsymbol{\psi}_j = \frac{m_j}{2} \int v_j^2 \mathbf{V}_j f_j d\mathbf{v}_j = \frac{1}{2} \rho_j m_j \overline{V_j^2 \mathbf{V}_j} = \mathbf{q}_j \quad (18-27)$$

the heat flux vector of the  $j$ th species.

It should be clear at this point that once we have an expression for  $f_j(\mathbf{r}, \mathbf{v}_j, t)$ , we can calculate all the fluxes and hence all the transport properties of a dilute gas. What we need now is  $f_j$ , or at least an equation that gives  $f_j$  as its solution. The only equation we have up to now is Eq. (18-14) with  $n = 1$ , and it can be seen that this also contains  $f_j^{(2)}$ . As we said earlier, nobody has found a successful way to uncouple this system. In the next section we shall derive an equation for  $f_j$ , the Boltzmann equation, which is the fundamental equation of the rigorous kinetic theory of gases.

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